

# Billiards with refraction: an example from Celestial Mechanics

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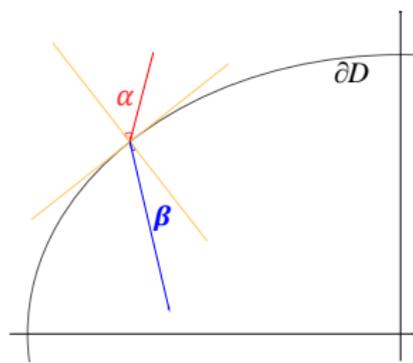
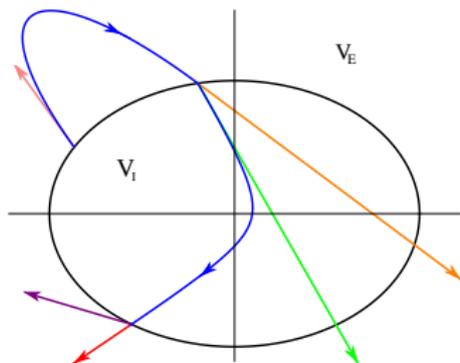
## Dynamical model

Let  $D \subset \mathbb{R}^2$  be a regular domain,  $0 \in D^\circ$ , and consider the orbits with zero energy subjected to the potential

$$V(z) = \begin{cases} V_I(z) = \mathcal{E} + h + \frac{\mu}{|z|} & \text{if } z \in D, \\ V_E(z) = \mathcal{E} - \frac{\omega^2}{2}|z|^2 & \text{if } z \notin \bar{D}, \end{cases}$$

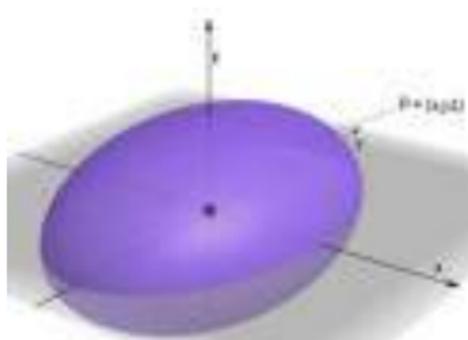
with  $\mathcal{E}, h, \mu, \omega > 0$ , while on the boundary  $\partial D$  the following junction rule (*Snell's law-type*) holds:

$$\sqrt{V_E(z)} \sin \alpha = \sqrt{V_I(z)} \sin \beta$$



# Motivations - Black Hole in an elliptical Galaxy

Suppose to have an ellipsoidal galaxy with constant density and a Black Hole at its center.



$$\begin{cases} V_{Gal}(P) = -\frac{\omega_x^2}{2}x^2 - \frac{\omega_y^2}{2}y^2 - \frac{\omega_z^2}{2}z^2 + C_G \\ V_{BH}(P) = \frac{\mu}{\sqrt{x^2+y^2+z^2}} + C_{BH} \end{cases}$$

$$\omega_x, \omega_y, \omega_z > 0, C_G, C_{BH} \in \mathbb{R}.$$

*BH's region of influence* =  $\tilde{D} = \{P \in \mathbb{R}^3 \mid |V_{BH}(P)| \gg |V_{Gal}(P)|\}$

$\pi_{xy}$  invariant under the dynamics: if  $\omega_x = \omega_y$  and  $0 < C_G < C_{BH}$ , setting  $D = \tilde{D} \cap \pi_{xy}$  we can use our 2-D model to study the motion of  $P$  on  $\pi_{xy}$ .

# First return map

$\partial D = \text{supp}(\gamma)$  with  $\gamma : \mathbb{R}/2\pi\mathbb{Z} \rightarrow \mathbb{R}^2$  regular closed curve; if  $(p_0, v_0) \in \partial D \times \mathbb{R}^2$  are the initial conditions of an outward-pointing orbit, there are  $\xi_0 \in \mathbb{R}/2\pi\mathbb{Z}, \alpha_0 \in [-\pi/2, \pi/2]$  such that, if  $\hat{t}(\xi), \hat{n}(\xi)$  are the tangent and inner normal unit vectors:

$$p_0 = \gamma(\xi_0), v_0 = \sqrt{2V_E(p_0)}(\sin \alpha_0 \hat{t}(\xi_0) + \cos \alpha_0 \hat{n}(\xi_0))$$

$\Rightarrow$  the pair  $(\xi_0, \alpha_0)$  completely determines the initial conditions on the boundary  $\partial D$ . We define the *first return map*

$$F : \mathbb{R}/2\pi\mathbb{Z} \times \left[-\frac{\pi}{2}, \frac{\pi}{2}\right] \rightarrow \mathbb{R}/2\pi\mathbb{Z} \times \left[-\frac{\pi}{2}, \frac{\pi}{2}\right]$$

$$(\xi_0, \alpha_0) \xrightarrow[\text{refraction E-I}]{\text{outer arc}} (\tilde{\xi}, \tilde{\alpha}) \xrightarrow[\text{refraction I-E}]{\text{inner arc}} (\xi_1, \alpha_1)$$

## Variational approach

Define the generating function

$$S : \mathbb{R}/2\pi\mathbb{Z} \times \mathbb{R}/2\pi\mathbb{Z} \rightarrow \mathbb{R} \quad S(\xi_0, \xi_1) = S_E(\xi_0, \tilde{\xi}) + S_I(\tilde{\xi}, \xi_1)$$

$$\begin{aligned} S_{E \setminus I}(a, b) &= d_{E \setminus I}(\gamma(a), \gamma(b)) = \\ &= \min \left\{ \int_0^1 |\dot{\lambda}(t)| \sqrt{V_{E \setminus I}(\lambda(t))} dt \mid \lambda(t) \text{ piecewise differentiable} \right. \\ &\quad \left. \text{s.t. } \lambda(0) = \gamma(a), \lambda(1) = \gamma(b) \right\} \end{aligned}$$

and  $\tilde{\xi}$  such that, fixed  $\xi_0, \xi_1$ ,  $\partial_{\tilde{\xi}}(S_E(\xi_0, \tilde{\xi}) + S_I(\tilde{\xi}, \xi_1))|_{\tilde{\xi}=\tilde{\xi}} = 0$ .

Conjugated actions:  $l_0 = -\partial_{\xi_0} S(\xi_0, \xi_1) = \sqrt{V_E(\gamma(\xi_0))} \sin \alpha_0$

$$l_1 = \partial_{\xi_1} S(\xi_0, \xi_1) = \sqrt{V_E(\gamma(\xi_1))} \sin \alpha_1$$

$$\begin{aligned} \mathcal{F} : \mathbb{R}/2\pi\mathbb{Z} \times \left( -\sqrt{\mathcal{E} - \frac{\omega^2}{2}}, \sqrt{\mathcal{E} - \frac{\omega^2}{2}} \right) &\rightarrow \mathbb{R}/2\pi\mathbb{Z} \times \left( -\sqrt{\mathcal{E} - \frac{\omega^2}{2}}, \sqrt{\mathcal{E} - \frac{\omega^2}{2}} \right) \\ (\xi_0, l_0) &\mapsto (\xi_1, l_1) \end{aligned}$$

When  $S(\xi_0, \xi_1)$  is well defined,  $\mathcal{F}$  is **conservative**.

## Circular case

If  $D$  is a disk of radius 1,  $\mathcal{F}$  is a **shift map**  
 $(\xi_0, l_0) \mapsto (\xi_1, l_1) = (\xi_0 + \theta(l_0), l_0)$ , where

$$\theta(l) = \begin{cases} \arctan\left(\frac{\varepsilon - 2l^2}{l\sqrt{4\varepsilon - 2(2l^2 + \omega^2)}}\right) + 2 \arccos\left(\frac{2l^2 - \mu}{\sqrt{4(\varepsilon + h)l^2 + \mu^2}}\right) - 2\pi & \text{if } l > 0 \\ 0 & \text{if } l = 0 \\ \arctan\left(\frac{\varepsilon - 2l^2}{l\sqrt{4\varepsilon - 2(2l^2 + \omega^2)}}\right) - 2 \arccos\left(\frac{2l^2 - \mu}{\sqrt{4(\varepsilon + h)l^2 + \mu^2}}\right) + \pi & \text{if } l < 0 \end{cases}$$

### Proposition

There is  $\tilde{\mathcal{I}} \subset \mathcal{I} = \left(-\sqrt{\varepsilon - \frac{\omega^2}{2}}, \sqrt{\varepsilon - \frac{\omega^2}{2}}\right)$ ,  $|\tilde{\mathcal{I}}| \leq 10$ , such that for every  $\xi_0 \in \mathbb{R}/2\pi\mathbb{Z}$ ,  $l_0 \in \mathcal{I} \setminus \tilde{\mathcal{I}}$ :

- $\mathcal{F}$  is well defined and conservative;
- $\mathcal{F}$  satisfies the **twist condition**

$$\frac{\partial \xi_1}{\partial l_0}(\xi_0, l_0) \neq 0.$$

# Orbits on the circle

Orbit:  $\{(\xi_k, l_k)\} = \{\mathcal{F}^k(\xi_0, l_0)\}_{k \in \mathbb{Z}}$ .

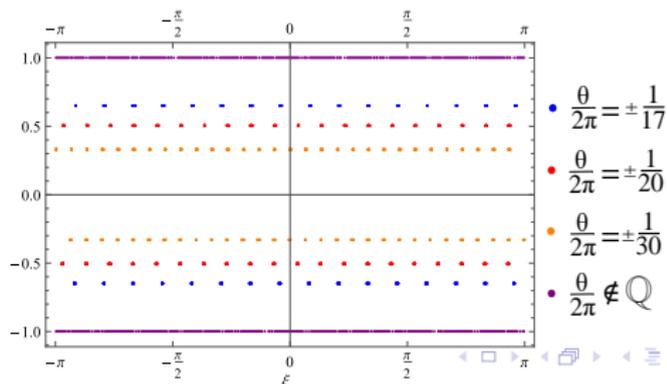
Rotation number:  $\rho(\xi_0, l_0) = \lim_{k \rightarrow \infty} \frac{\xi_k}{k} = \theta(l_0)$  (circular case).

If  $\theta(l_0) = 2\pi \frac{p}{q}$ ,  $p, q \in \mathbb{Z} \implies (\xi_q, l_q) = (\xi_0 + 2\pi p, l_0) \equiv (\xi_0, l_0)$   
 $\implies (\xi_0, l_0)$  is  $(p, q)$  - periodic.

## Theorem

$\exists C = C(\mathcal{E}, \omega, h, \mu) > 0$  such that for every  $\rho \in (-C, C)$  except for a finite number of values  $\exists l_0^\pm \in \mathcal{I} \setminus \tilde{\mathcal{I}}$  s.t. for every  $\xi_0 \in \mathbb{R}/2\pi\mathbb{Z}$   
 $\rho(\xi_0, l_0^\pm) = \rho$ .

$\mathcal{E} = 7, \omega^2 = 3$   
 $h = 2, \mu = 15$   
 $C \approx 0.391$



## General case: homotetic fixed point

Suppose that  $\gamma : \mathbb{R}/2\pi\mathbb{Z} \rightarrow \mathbb{R}^2$  is a regular closed curve,  $\text{supp}(\gamma) = \partial D$ ,  $0 \in D^\circ$ .

### Proposition

If  $\bar{\xi} \in \mathbb{R}/2\pi\mathbb{Z}$  is such that  $\dot{\gamma}(\bar{\xi}) \perp \gamma(\bar{\xi}) \Rightarrow (\bar{\xi}, 0)$  is a fixed point for  $\mathcal{F}$ , corresponding to the homotetic solution of initial conditions

$$(p_0, v_0) = \left( \gamma(\bar{\xi}), \frac{\sqrt{2V_E(\gamma(\bar{\xi}))}\gamma(\bar{\xi})}{|\dot{\gamma}(\bar{\xi})|} \right)$$

The linear stability of  $(\bar{\xi}, 0)$  can be studied by computing the Jacobian matrix of  $\mathcal{F}$ .

*Variational methods + Levi-Civita regularization + Implicit function theorem*

↓

$$D\mathcal{F}(\bar{\xi}, 0) = \begin{pmatrix} \frac{\partial \xi_1}{\partial \xi_0}(\bar{\xi}, 0) & \frac{\partial \xi_1}{\partial I_0}(\bar{\xi}, 0) \\ \frac{\partial I_1}{\partial \xi_0}(\bar{\xi}, 0) & \frac{\partial I_1}{\partial I_0}(\bar{\xi}, 0) \end{pmatrix} \rightsquigarrow \lambda_1, \lambda_2 \text{ eigenvalues}$$

## Homotetic fixed points - Linear stability

### Proposition

Let  $p(\lambda)$  be the characteristic polynomial of  $D\mathcal{F}(\bar{\xi}, 0)$ , and  $\Delta$  its discriminant.

Since  $\det(D\mathcal{F}(\bar{\xi}, 0)) = 1$ , then:

- $\Delta > 0 \Rightarrow \lambda_1, \lambda_2 \in \mathbb{R}, \lambda_2 = 1/\lambda_1, \lambda_1 \neq \lambda_2 \Rightarrow |\lambda_1| > 1, |\lambda_2| < 1$ , then  $(\bar{\xi}, 0)$  is an unstable saddle point;
- $\Delta < 0 \Rightarrow \lambda_1, \lambda_2 \in \mathbb{C} \setminus \mathbb{R}, \lambda_1 = \bar{\lambda}_2, \lambda_1 \neq \lambda_2 \Rightarrow |\lambda_1| = |\lambda_2| = 1$  and  $(\bar{\xi}, 0)$  is a stable center point;
- $\Delta = 0 \rightarrow$  degenerate case (e.g. circular domain).

The discriminant  $\Delta$  depends on the physical parameters  $\mathcal{E}, h, \mu, \omega$  and on  $\gamma$  and  $\bar{\xi}$  through  $|\gamma(\bar{\xi})|, |\dot{\gamma}(\bar{\xi})|$  and  $k(\bar{\xi})$ , i.e. the curvature of  $\gamma$  in  $\bar{\xi}$ : fixed  $\bar{\xi}$  such that  $(\bar{\xi}, 0)$  is a homotetic fixed point, changing the values of the physical parameters may modify its stability properties  $\rightarrow$  **bifurcations**.

## Elliptic case - Bifurcation for $\mu$

Suppose that  $\partial D$  is an ellipse parametrized by  $\gamma(\xi) = (a \cos \xi, b \sin \xi)$ ,  $a = 1$ ,  $b = a\sqrt{1 - e^2}$ ,  $0 \leq e < 1 \Rightarrow (0, 0)$  and  $(\pi/2, 0)$  are fixed points for  $\mathcal{F}$ .

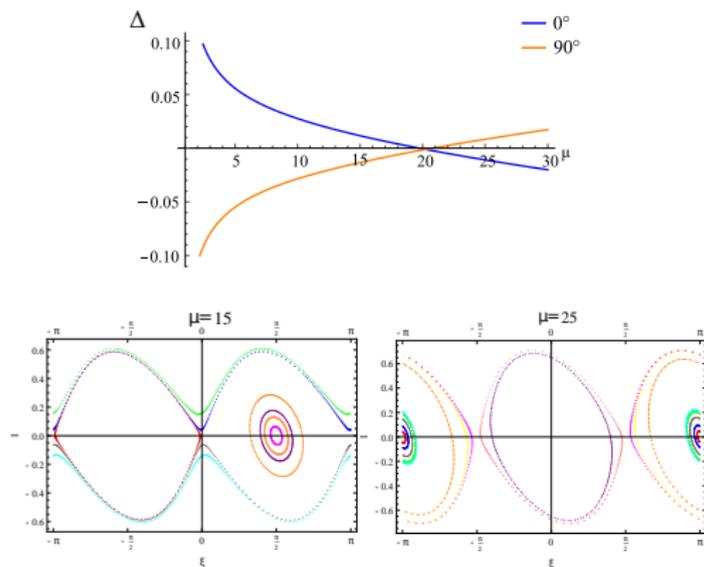
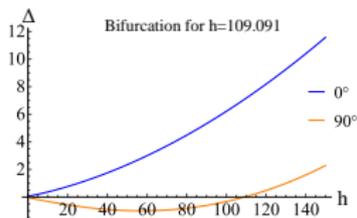


Figure:  $\mathcal{E} = 2.5$ ,  $h = 2$ ,  $\omega = \sqrt{2}$ ,  $e = 0.1$ . Up: value of  $\Delta$  for  $(0, 0)$  and  $(\pi/2, 0)$  and  $\mu \in [1, 30]$ . Down: Poincaré sections for different values of  $\mu$ , 200 points

# Elliptic case - Bifurcation for $h$



The stability behaviour of the equilibria suggests the presence of a pitchfork bifurcation, with the arising of a new equilibrium point between 0 and  $\pi/2$ .

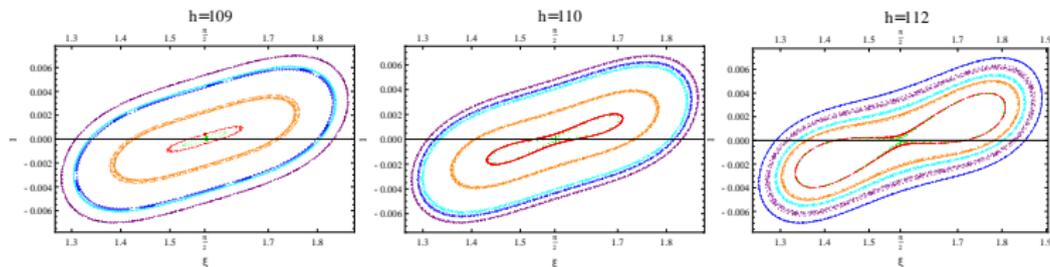
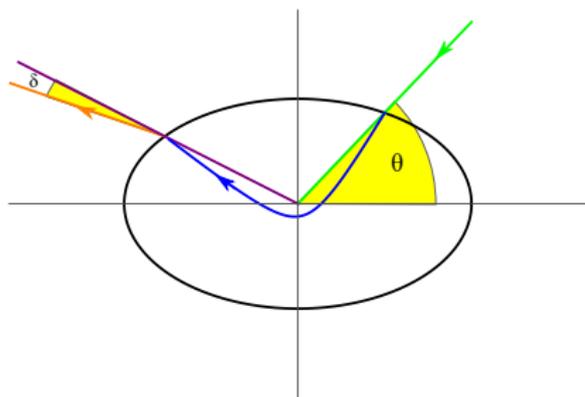
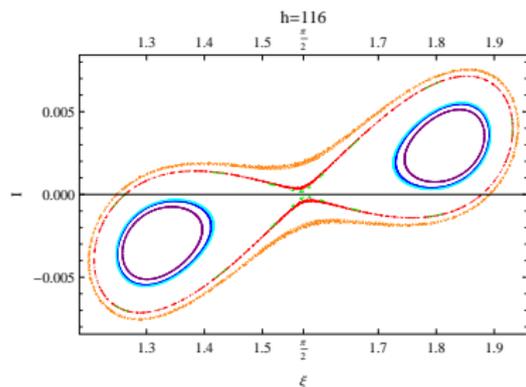


Figure:  $\mathcal{E} = 2.5, \mu = 2, \omega = \sqrt{2}, e = 0.1$ . Up: value of  $\Delta$  for  $(0,0)$  and  $(\pi/2,0)$  and  $h \in [0, 150]$ . Down: Poincaré sections for different values of  $h$ , 2000 points.

## Elliptic case - Search for new equilibria



If  $\lambda_{1/2}^{(k)}$  are the eigenvalues of  $D\mathcal{F}^k(\bar{\xi}, 0) \Rightarrow \lambda_{1/2}^{(k)} = \lambda_{1/2}^k \Rightarrow$  the bifurcation values are the same for *all the iterates*.

The Poincaré sections show the arising of an orbit of minimal period 2 for  $h >$  bifurcation value.

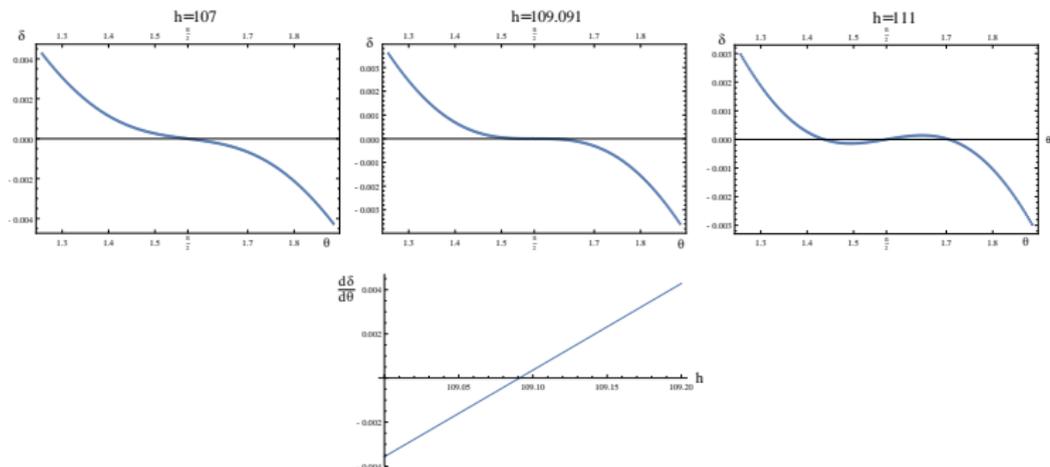
*Idea:* search for 2-periodic **brake orbits** via the **shooting method**.

We search for 2-periodic orbits which are homotetic in their outer arcs: consider the free fall map

$$\Phi : [0, 2\pi] \rightarrow [0, 2\pi], \quad \theta \mapsto \delta$$

## Elliptic case - Search of new equilibria

If  $\delta = \Phi(\theta) = 0$ , the outer branches are both homotetic, and the whole orbit is 2-periodic.



Up: plot of  $\delta = \Phi(\theta)$  in a neighborhood of  $\pi/2$  for different values of  $h$ .  
Down: first derivative of  $\delta^{\pi/2} = \Phi(\pi/2)$  as a function of  $h$ .

**Conclusion:** the 2-periodic orbits which appear when  $h >$  bifurcation value are brake.

## Further research

- *Elliptic case*: systematic study of the orbits near to the homotetics for every value of the eccentricity;
- *General case*: perturbative methods on the circular case, with a small deformation of the boundary  $\partial D$ . In particular:
  - **KAM theory**: existence of quasi-periodic invariant tori with diophantine rotation number for the perturbed system;
  - **Mather theory** (Poincaré Birkhoff theorem): existence of invariant orbits with prescribed rotation number, under the hypothesis of *twist condition*.

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